

Forward and Inverse Analyses of Smart Compliant Mechanisms for Path Generation

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Abstract

In this paper, an open loop compliant mechanism consisting of two elastic links actuated through piezoelectric actuators has been analyzed. The links may be joined through rigid or elastic hinge connection. The effect of piezoelectric actuators placed on an elastic beam is considered as two concentrated self balancing moments acting on the beam at the edges of the actuator. Considering the large deflection of the two links of the mechanism under these self balancing moments as well as end loads, the forward and inverse analyses of the mechanism are carried out. Two numerical schemes, namely, non-linear shooting and Adomian decomposition methods have been used for solving these problems. Numerical results are included to demonstrate the potential of the proposed methods.

Keywords: Complaint Mechanism, Smart Actuators

1. Introduction

Compliant mechanisms are composed of elastic links whose deformations are utilized to produce the desired output motion for a given input actuation. Since these are joint-less and monolithic in nature, there is no friction and wear; resulting in improved repeatability, reduced maintenance and simpler manufacturing [1]. But in comparison to a rigid body mechanism, a compliant mechanism yields smaller workspace. Furthermore, large deformation of its links enforces the consideration of geometric and material non-linearity, thus making its design and synthesis a challenging task. Towards the synthesis of compliant mechanisms, two complementary methods are found in the literature: the pseudo-rigid body model [2, 3], which utilizes the synthesis technique of rigid body mechanisms and the topology synthesis method [4], which uses the structural optimization technique. All these methods are used to design compliant mechanisms for function or path generation tasks. The mechanisms thus produced are mainly of *passive* type, i.e., the sources of input forces and/or moments are not considered in the design. In order to complete the design, proper selection of the actuators is necessary. In this respect, use of smart material actuators is best suited for the actuation of compliant mechanisms because of their wide variety and capability. The combined system of smart material actuators and compliant mechanisms has certain advantages as reported in [5]. These smart material actuated compliant mechanisms can be called *active*

compliant mechanisms. But the design and analyses of the resulting smart material actuated compliant mechanisms involve numerous complexities. Besides the requirement of the considerations of geometric and material non-linearity, the input characteristics of the compliant mechanism and the output characteristics of the smart material actuators should also be matched. A theoretical study of such a smart material actuator driven compliant mechanism with the objective of path generation is presented here.

In the literature of compliant mechanisms, mostly closed chain deformable structures are designed or synthesized. The term closed chain mechanism bears the same meaning as in the traditional rigid body mechanism. In this paper, the forward and inverse analyses of an open chain compliant mechanism, consisting of two deformable links joined through a rigid joint or an elastic hinge and actuated through smart material actuators, have been presented. This mechanism is very similar to a traditional open chain manipulator consisting of rigid links and kinematic pairs, actuated through electric motors or pneumatic drives at each joint. In the present case, the links are elastic and each of them is independently actuated through piezoelectric actuators. The effect of two piezoelectric actuators, mounted on the two opposite sides of an elastic beam and actuated through out of phase voltage, is simulated as two self balancing moments (equal in magnitude and opposite in sense) acting just at the edge of the actuators as reported in [6]. The relation between the applied input voltage to piezoelectric actuators and the output moments exerted by these on the beam is explicitly derived in [6, 7]. In case of links actuated through multiple actuators, each of the them embedded on the surfaces of the beam is replaced by a pair of self balancing moments, the magnitude of which depend on the applied voltage across the actuator and its material properties. This idea of replacing the piezoelectric actuators with pairs of self balancing moments has been used in [8]. But in that case, the linear beam theory is used to determine the static deflection. The concept of constructing a flexible mechanism consisting of multiple deformable links, connected in series with rigid or elastic connection and driven through smart material actuators embedded in each of the links, is not found in the literature. Furthermore, there is always an end force and moment acting on the mechanism simulating the payload on the same. This type of flexible mechanisms may be used where precise and repeatable motion is required, such as in cell harvesting.

The present paper is organized in the following manner. First the equilibrium equation of a flexural beam is obtained using Euler-Bernoulli beam theory and the resulting non-linear differential equation is solved using

Shooting and Adomian decomposition techniques. Secondly, the forward analysis (i.e., the determination of the end point trajectory given the actuating moments) has been carried out for the two link flexible mechanism, where the links can be joined through either a rigid or an elastic connection. The actuating moments are distributed over each of the compliant links and for a given range of moments, the workspace has been generated. Finally, the inverse problem (i.e., to determine the moments required to reach the desired points for a given trajectory within the workspace) has been solved.

It should be mentioned that the complete solution of a path generation problem by a smart compliant mechanism must include the dynamics of the actuator and that of the mechanism. However, for low speed actuation and actuators with fast dynamics (e.g., piezoelectric and magnetostrictive actuators) the static analysis discussed in this paper can provide a useful first approximation.

2. Cantilever beam subjected to self-balanced moment and external load

In this section, the large deflection of a cantilever beam under self-balanced moments as well as external forces has been formulated and solved.

2.1. Non-linear shooting method for a single link

Figure (1) shows the deformed shape of a cantilever beam, the model of a compliant segment, subjected to two equal and opposite moments applied at intermediate locations. Furthermore, there is a non-following end force and a moment. The intermediate moments are acting at distances l_1 and l_2 from the fixed end along the length of the deformed beam. The resulting bending moment at any point (x, y) on the beam is given by

$$M_{(x,y)} = P(a-x) + nP(b-y) + M_p [u(s-l_1) - u(s-l_2)] + M_e \quad (1)$$

where s is the distance of the point from the fixed end along the length of the beam, P and nP are the vertical and horizontal components of the non-following end force F , M_p is the moment applied by the actuators, $u(s)$ is the unit step function defined as $u(s) = 0$ for $s < 0$ and $u(s) = 1$ for $s \geq 0$ and M_e is the external moment applied at the free end. From Eq. (1) and the Euler-Bernoulli moment-curvature relationship one gets

$$EI \frac{d\theta}{ds} = P(a-x) + nP(b-y) + M_p [u(s-l_1) - u(s-l_2)] + M_e \quad (2)$$

where $\frac{d\theta}{ds}$ denotes the curvature of the beam and EI is its flexural rigidity.

Differentiating Eq. (2) with respect to \bar{s} and substituting $\frac{dx}{ds} = \cos \theta$ and $\frac{dy}{ds} = \sin \theta$ one gets

$$\left. \begin{aligned} \frac{d^2\theta}{d\bar{s}^2} &= -\frac{PL^2}{EI} (\cos \theta + \sin \theta) \\ &+ \frac{M_p L}{EI} [\delta(\bar{s} - \bar{l}_1) - \delta(\bar{s} - \bar{l}_2)] \end{aligned} \right\} \quad (3)$$

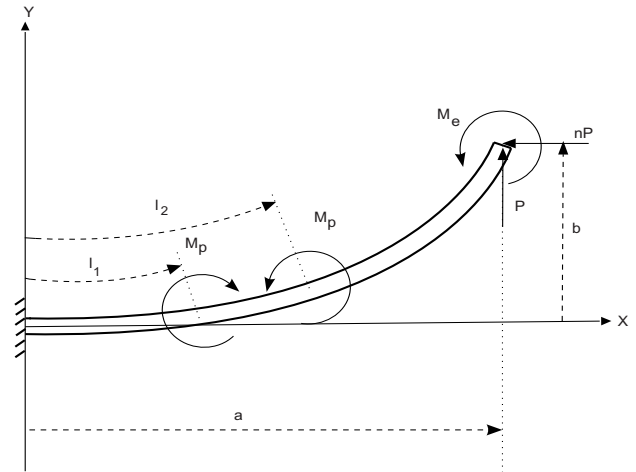


Figure 1: Deformed shape of an elastic link

where $\bar{s} = \frac{s}{L}$, $\bar{l}_1 = \frac{l_1}{L}$, $\bar{l}_2 = \frac{l_2}{L}$, L is the length of the beam and $\delta(s)$ is the Dirac-Delta function defined as $\delta(s) = 0$ for $s \neq 0$ and $\delta(s) \rightarrow \infty$ at $s = 0$. Introducing the non-dimensional load parameter $\alpha = \frac{PL^2}{EI}$ and the non-dimensional moment parameter $\kappa = \frac{M_p L}{EI}$ in Eq. (3) the boundary value problem is posed as

$$\left. \begin{aligned} D.E. \quad \frac{d^2\theta}{d\bar{s}^2} &= -\alpha (\cos \theta + n \sin \theta) + \\ &\kappa [\delta(\bar{s} - \bar{l}_1) - \delta(\bar{s} - \bar{l}_2)] \\ B.C. \quad \theta|_{\bar{s}=0} &= 0, \quad \frac{d\theta}{d\bar{s}}|_{\bar{s}=1} = \tau \end{aligned} \right\} \quad (4)$$

where $\tau = \frac{M_e L}{EI}$ is the non-dimensional external moment acting at the free end of the beam.

In the non-linear shooting method, the boundary value problem (B.V.P) is converted into an initial value problem (I.V.P) with an assumed curvature at the fixed end i.e., $c = \frac{d\theta}{d\bar{s}}|_{\bar{s}=0}$. First the differential equation is solved using fourth order Runge-Kutta method with an assumed value of c . Then the value of c is modified till the second boundary condition is satisfied. Here I.V.P is posed as

$$\left. \begin{aligned} D.E. \quad & \frac{d^2\theta}{d\bar{s}^2} = -\alpha(\cos\theta + n\sin\theta) + \\ & \kappa[\delta(\bar{s} - \bar{l}_1) - \delta(\bar{s} - \bar{l}_2)] \\ B.C. \quad & \theta|_{\bar{s}=0} = 0, \quad \left. \frac{d\theta}{d\bar{s}} \right|_{\bar{s}=0} = m_k \end{aligned} \right\} \quad (5)$$

where m_k is assumed to be the curvature at the fixed end at the k^{th} iteration step. Thus, the error involved is determined as $error = \left[\left(\frac{d\theta}{d\bar{s}} \right)_{\bar{s}=1} - \tau \right]$ and this is to be made zero by properly modifying m_k . In this paper, Newton-Raphson method has been used. Now m_k in the k^{th} step is calculated from that of the $(k-1)^{\text{th}}$ step using

$$m_k = m_{k-1} - \frac{(error)}{\frac{\partial}{\partial m} \left[\frac{d\theta}{d\bar{s}} \right]_{\bar{s}=1}}. \quad (6)$$

Thus, at each step m_k is modified using Eq. (6) till the error is less than some allowable tolerance and finally the solution is achieved. The proof of convergence of this method is explicitly shown in [9]. Obtaining the solution for $\theta(\bar{s})$, \bar{x} and \bar{y} is obtained using $\frac{d\bar{x}}{d\bar{s}} = \cos\theta$ and $\frac{d\bar{y}}{d\bar{s}} = \sin\theta$ respectively.

For multiple patches, Eq. (4) can be modified as

$$\left. \begin{aligned} D.E. \quad & \frac{d^2\theta}{d\bar{s}^2} = -\alpha(\cos\theta + n\sin\theta) + \\ & \sum_i \kappa_i [\delta(\bar{s} - \bar{l}_1^i) - \delta(\bar{s} - \bar{l}_2^i)] \\ B.C. \quad & \theta|_{\bar{s}=0} = 0, \quad \left. \frac{d\theta}{d\bar{s}} \right|_{\bar{s}=1} = \tau \end{aligned} \right\} \quad (7)$$

where i denotes the number of piezo patches and \bar{l}_1^i and \bar{l}_2^i are the distances of the two ends of the i^{th} smart actuator from the fixed end of the beam. Eq. (7) is used to solve the large deflection beam problem subjected to multiple patches.

Thus, the deformed shape of a single piece flexible link subjected to multiple pairs of self-balancing moments and end loadings (forces and moments) is obtained using the shooting technique. This method will be iteratively used in modeling the two link flexible mechanism.

2.2. Adomian decomposition method for a single link

Adomian decomposition method provides an analytical solution to non-linear ordinary and partial differential equations in series form. This method has been used to solve numerous engineering and scientific problems governed by non-linear ordinary and partial differential equations [10]. The non-linear terms in the governing differential equation are decomposed into polynomials

using Taylor's series, known as Adomian polynomials, and are used to obtain the analytical solution. The method and its proof of convergence have been elaborately explained in [11].

Our aim is to solve the boundary value problem governed by the non-linear differential equation given by Eq. (4). Integrating Eq. (4) twice with respect to \bar{s} , using the B.C.s' and upon simplification, the governing equation takes the form

$$\left. \begin{aligned} \theta_1(\bar{s}) &= c\bar{s} + \alpha \int_0^{\bar{s}} \int_0^t (\cos\theta + n\sin\theta) d\bar{s} dt \\ & \quad \text{for } 0 \leq \bar{s} < \bar{l}_1 \\ \theta_2(\bar{s}) &= c\bar{s} + \alpha \int_0^{\bar{s}} \int_0^t (\cos\theta + n\sin\theta) d\bar{s} dt + \kappa(\bar{s} - \bar{l}_1) \\ & \quad \text{for } \bar{l}_1 \leq \bar{s} < \bar{l}_2 \\ \theta_3(\bar{s}) &= c\bar{s} + \alpha \int_0^{\bar{s}} \int_0^t (\cos\theta + n\sin\theta) d\bar{s} dt + \kappa(\bar{l}_2 - \bar{l}_1) \\ & \quad \text{for } \bar{l}_2 \leq \bar{s} < 1 \end{aligned} \right\} \quad (8)$$

where $c = \left. \frac{d\theta}{d\bar{s}} \right|_{\bar{s}=0}$, $\theta_1(\bar{s})$, $\theta_2(\bar{s})$ and $\theta_3(\bar{s})$ are

respectively the angle that the tangent at any point on the three segments of the beam makes with the X axis. Each of the integral equations is solved by expanding the non-linear term within the integral using Adomian polynomials about $\theta_{\bar{s}=0} = 0$. The procedures to obtain the coefficients of Adomian polynomials corresponding to a particular non-linear term are explicitly explained in [12].

Unlike the shooting method, here the unknown c is determined by satisfying the moment boundary condition specified at the free end i.e., $\left. \frac{d\theta}{d\bar{s}} \right|_{\bar{s}=1} = \tau$. Once $\theta(\bar{s})$ is obtained for each of the segments, $(\bar{x}(\bar{s}), \bar{y}(\bar{s}))$ can be computed using $\frac{d\bar{x}}{d\bar{s}} = \cos\theta$ and $\frac{d\bar{y}}{d\bar{s}} = \sin\theta$. This can be done either by expanding the sine and cosine terms in Taylor's series or by numerical integration. Since the Adomian polynomial used in approximating the non-linear term in Eq. (8) is truncated, the corresponding solution is not exact and thus c^0 continuity will not be maintained in between the segments. In order to enforce the continuity, the errors i.e., $(\theta_1 - \theta_2)$ at $\bar{s} = \bar{l}_1$ and $(\theta_2 - \theta_3)$ at $\bar{s} = \bar{l}_2$ are deducted from $\theta_2(\bar{s})$ and $\theta_3(\bar{s})$, respectively. Thus the entire beam configuration is obtained.

This procedure can be extended for a beam with multiple patches also, where the number of equations increases accordingly. For each additional actuator two new equations are to be solved.

Thus, Adomian decomposition method can be used to obtain the deformed shape of a single piece flexible link subjected to multiple pairs of self-balancing moments as well as end loadings. In the next section, this method has been iteratively used to solve the two link mechanism problem.

3. Forward analysis of a two link mechanism

The objective of the forward analysis is to determine the position of the end point of the mechanism for given actuating moments M_p^i applied by the actuators while manipulating an external load at the end.

3.1 Non-linear shooting method

The method elucidated in section 2.1 for a single elastic link and multiple patches is iteratively used to solve the two link mechanism problem. The equilibrium of each link is considered separately under all the external forces and moments as well as the mutual reactions between the neighboring links. The algorithm is shown step by step in the next section. All the symbols have been defined in previous sections.

3.1.1 Two link flexible mechanism with rigid connection

Figure (2) shows a flexible mechanism consisting of two elastic links connected rigidly. The angle between the links is β . Each of the links is of length L and actuated through a pair of self-balancing moments M_p^1 and M_p^2 , respectively. The components of the non-following force P and nP and the moment at the free end M_e . The free body diagram of each of the elastic links is shown in Fig. (3). Here M_x is the reaction moment acting between the links. Two coordinate systems are depicted in this figure. The coordinate system (X, Y) is the fixed reference, with its abscissa along the undeflected axis of the link-I. The final coordinates of the end points (X_1, Y_1) and (X_2, Y_2) of link-I and link-II, respectively, have been defined in this coordinate system. The (x, y) coordinate system with its origin at the end of link-I and its abscissa along the undeflected axis of the link-II has been used to write the equilibrium equation of the second link. The algorithm for solving the forward analysis is as follows.

Step I: The normalized end force α and moment τ are determined for a given payload using the normalization scheme described above.

Step II: Configuration of the link-I is obtained using α , n , κ_1 and τ_x ; where $\kappa_1 = \frac{M_p^1 L}{EI}$, the normalized actuating moment on the link-I and $\tau_x = \frac{M_x L}{EI}$, the normalized end moment, which is initially assumed to be zero.

Step III: Configuration of link-II is determined corresponding to α' , n' , κ_2 and τ , where α' and n' are obtained by transforming the end forces in the coordinate system (x, y) (obtained through rotation of the original reference (X, Y) by an angle equal to $(\theta_e - \beta)$); where

θ_e is the angle that the tangent drawn at the end of link-I makes with the X axis, $\kappa_2 = \frac{M_p^2 L}{EI}$ is the normalized actuating moment on link-II and $\tau = \frac{M_e L}{EI}$ is the normalized end moment. Since the joint is rigid, angle β remains unchanged even after deformation of the links.

Step IV: From the final configuration of the link-II, the normalized moment τ_x at the end of link-I which is equal and opposite to the moment at the beginning of link-II is computed using

$$\tau_x = \alpha(\bar{X}_2 - \bar{X}_1) + n\alpha(\bar{Y}_2 - \bar{Y}_1). \quad (9)$$

where $\bar{X} = \frac{X}{L}$ and $\bar{Y} = \frac{Y}{L}$.

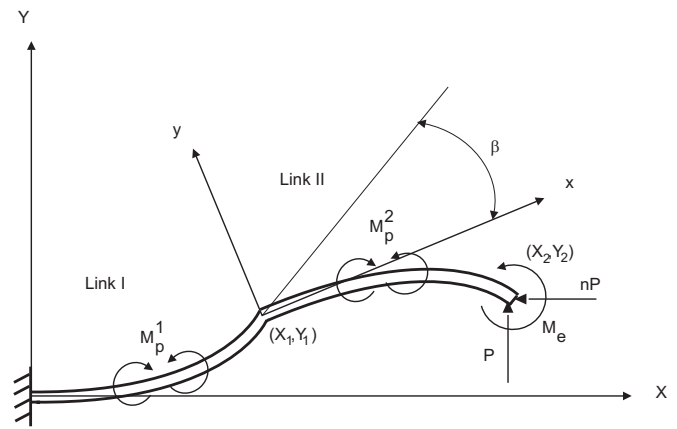


Figure 2: Deformed two-link flexible mechanism

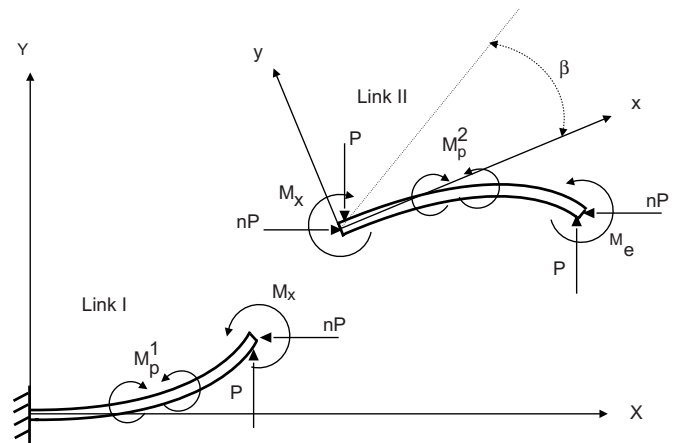


Figure 3: Free body diagram of two elastic links connected through rigid joints

Step V: The non-zero value of τ_x , the non-dimensional moment applied by link-II on link-I, yield a new configuration of link-I and simultaneously changes the position and orientation of the (x, y) axes and thus yields new values of α' and n' . In this way a new configuration of link-II is obtained.

Step VI: Again the value of the normalized reaction moment τ_x is updated and the deformed configuration for link-I is computed, resulting in a new value of θ_e .

Step VII: The procedures described in step III to step VI are continued until and unless the difference in the old and the new value of τ_x is less than the allowable tolerance.

The convergence in the value of τ_x means both the links are in equilibrium under the action of the actuating moments κ_1 and κ_2 as well as the end force α and the end moment τ .

3.1.2 Two link flexible mechanism with elastic hinge connection

The free body diagrams of the two links as well as that of the hinge are shown in Fig. (4). In the undeformed condition of the hinge, the angle between the links is β , and ψ is the change in the value of β i.e., the angular deformation of the hinge. The coordinate systems and the rest of the symbols remain same as defined in the previous section. In this case, apart from the equilibrium of the links, the equilibrium of the hinge is also to be considered. The steps of the solution algorithm are detailed below.

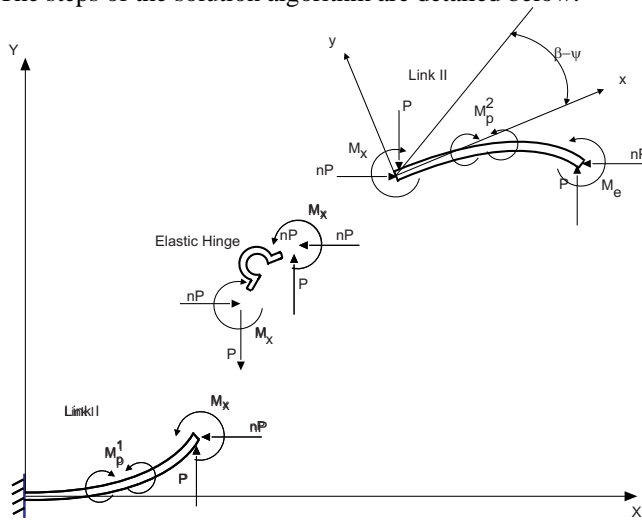


Figure 4: Free body diagram of two elastic links and the elastic hinge

Steps I and II remain the same as in section 3.1.1. In step III, the end forces (α', n') are obtained by rotating the coordinate (X, Y) by an angle $(\theta_e - \beta - \psi)$ instead of $(\theta_e - \beta)$, where ψ (equal to the change in β) is the angular deflection of the elastic hinge connecting the two links. Following step-IV described in the previous section the normalized end moment τ_x is calculated. Since τ_x is known, ψ , the angular deflection of the elastic hinge is determined using the relation $\tau_x = \zeta\psi$, where ζ is the normalized stiffness of the elastic hinge. If K_ψ is the angular stiffness of the hinge, the normalized stiffness is obtained using $\zeta = \frac{K_\psi L}{EI}$, where EI is the flexural rigidity

of the links, each of length L . It may be considered as the relative stiffness of the hinge with respect to the links. In each iteration, the value of ψ is computed and incorporated for the transformation of the end forces from (X, Y) coordinate system to (x, y) coordinate system. Rests of the steps are same as in the case of elastic links with rigid connection reported in section 3.1.1.

3.2 Adomian decomposition method

The Adomian decomposition method discussed in section 2.2 provides a closed form solution of Eq. (8) i.e., $\theta(\bar{s})$ of a single link. Those solutions are used iteratively to solve the two link mechanism problem just as in the case of shooting method. The details of the steps are listed below.

3.2.1 Two link flexible mechanism with rigid connection

Step I: Using the expression of $\theta(\bar{s})$ for link-I obtained in section 2.2, the unknown $c = \frac{d\theta}{ds}\bigg|_{\bar{s}=0}$ is determined with the initial assumption of $\tau_x = 0$. Thus the deformed shape of link-I and hence θ_e is determined.

Step II: The end force parameters (α', n') are determined by transforming (α, n) from (X, Y) coordinate system to (x, y) coordinate system. The transformation involves the rotation of (X, Y) axes by an amount $(\theta_e - \beta)$. Since the connection is rigid, β remains unchanged.

Step III: The values of (α', n') and the end moment τ are utilized to compute the second link configuration. In this case, c for the link-II is determined using the expressions for $\theta(\bar{s})$ and $\frac{d\theta}{ds}\bigg|_{\bar{s}=1} = \tau$.

Step IV: Resulting (X_2, Y_2) lead to a non zero value of τ_x , calculated using Eq. (9). Consequently the end slope of link-I and in turn the values of (X_2, Y_2) are changed.

Step V: Steps I-IV are continued till the difference in the old and new τ_x is less than an allowable value.

3.2.2 Two link flexible mechanism with elastic hinge connection

In this case, all the steps except third one discussed in section 3.2.1 remain same. Since the hinge is deformable, the angle β between the two elastic links no longer remains unaltered. Thus, the coordinate transformation of the end forces from (X, Y) to (x, y) is obtained through the rotation of the reference coordinate (X, Y) by an angle $(\theta_e - \beta - \psi)$ as is done in section 3.1.2.

4. Inverse analysis of a two link mechanism

The objective of the inverse analysis is to determine the actuating moments M_p required in each of the links to guide the end point of the mechanism through some desired locations. Though both the methods described above are capable of handling multiple actuators, only one actuator at each of the link is considered. Furthermore, there are non-following end forces and moment acting at the end of the mechanism.

Both the shooting and Adomian decomposition methods can be used in this analysis. The latter method becomes computationally intensive with increasing load parameters. But the effectiveness and the accuracy of the former remains independent of the magnitude of the force parameters. Hence in this work, the shooting method is used for the inverse analysis. But for low values of load parameters, Adomian method may be advantageous. In any case, the idea of solving the inverse problem is same for both the methods. For a given end loading the workspace increases considerably with an elastic connection between the links. In order to reduce the size of the workspace, only the rigid connection between the links has been considered for the inverse problem.

4.1 Inverse analysis using shooting method

As discussed in section 3.1.1, the end point (X_2, Y_2) of the two link mechanism with rigid connection is a function of the self-balanced moments (M_p^1, M_p^2) and the end forces and moments (P, n, M_e) . This can be represented as

$$\left. \begin{aligned} X_2 &= \phi_1(M_p^1, M_p^2, P, n, M_e) \\ Y_2 &= \phi_2(M_p^1, M_p^2, P, n, M_e) \end{aligned} \right\} \quad (10)$$

While using one actuator in each of the links, two unknowns M_p^1 and M_p^2 can be solved from the two equations given by Eq. (10), because all the other variables are known. Hence the system is consistent and this is one of the reasons for choosing one actuator in each link.

Use of multiple actuators will definitely increase the size of the workspace as well as the mechanisms capability to trace more and more complex trajectories. But in that case, the number of unknowns is greater than the number of equations and thus optimization schemes might result in multiple solutions.

5. Results and Discussions

Figure (5a) shows the deformed cantilever beam configurations (actuated through two piezoelectric actuators) obtained by using both non-linear shooting and Adomian decomposition methods. The centers of the piezoelectric actuators are assumed to be at a distance of 25% and 45% of the beam length measured from the fixed

end. The normalized actuating moments $\kappa_i s'$ are kept constant while the end load parameters are varied. In this figure the occurrence of inflection points, depending on the combinations of α and τ , may be noted. Thus it is confirmed that both these methods, unlike the elliptic integral method [2], can handle inflection point without any special treatment.

Figure (5b) shows the path generated by the free end (X_j, Y_j) of the cantilever, as each actuator, placed at different locations, applies a moment κ varying from -1 to +1 in steps of 0.1. It is clearly seen that, by using a single link, the shape of the end point trajectory does not significantly alter by changing the location of the actuator. Furthermore, the movement along the X-axis is also insignificant. This calls for the necessity of using a two link mechanism to generate a complicated trajectory.

Figure (6a) shows the deformed configurations of a two link mechanism connected through a rigid joint for constant actuating moments with varied end conditions. Both Adomian and Shooting methods yield the same results. In this case the angle (β) between the two links is taken as 90° . Here, each of the links is embedded with one piezo actuator, i.e., two self-balancing moments are acting on each link.

Figure (6b) shows the deformed configurations of the mechanism with elastic hinge connection for various loading combinations. Comparing Figs. (6a) and (6b) it can be observed that, the mechanism consisting of an elastic hinge ($\zeta = 1$, as in Fig. (6b)) undergoes larger deflection than that with a rigid joint ($\zeta = \infty$, as in Fig. (6a)). In order to see the effect of hinge stiffness, the mechanism configurations are plotted in Fig (6c) for various stiffness values of the hinge. These results are computed using only the shooting method. Obviously the results tend to that with a rigid joint, as ζ increases.

Every mechanism has its own working area, within which it is designed to manipulate the payload. In case of path generation with precision points approach, the end point of the mechanism is desired to go through the accuracy points located on the prescribed path. These accuracy points are chosen in advance to generate the prescribed path. In order to have an idea of all these, the workspace of the mechanism can be generated with the full capacity of the actuators as well as the payload for which it is to be designed.

As mentioned earlier, only the non-linear shooting method is used for inverse analysis, because of its effectiveness in computation. Figure (7a) shows the generation of a straight line with seven accuracy points. First, the actuating moments are computed by the inverse analysis for a given accuracy point. Thereafter, the actuating moments are used in the forward analysis to obtain the points actually reached. The maximum error, defined as the distance between the target accuracy point and the point reached, is in the order of 10^{-6} times the link length. The actuating moments $\kappa_i s'$ required to generate the straight line in Fig. (7a) are shown in Fig. (7b). This shows that a gradual change in moments (i.e., the voltage applied to the piezo) is required for each of the

piezoelectric actuators. Similar results for generating a figure of '8' are shown in Fig. (7c) - (7d).

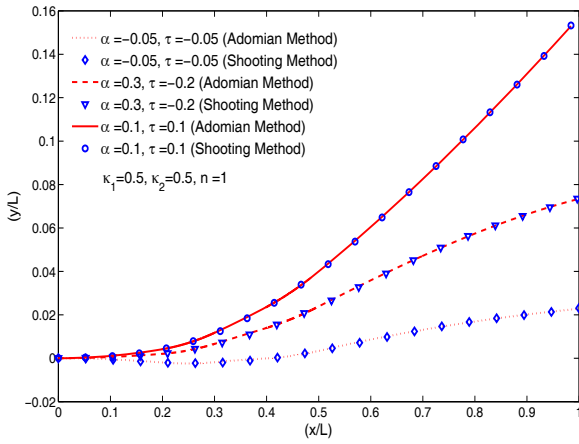


Figure 5a

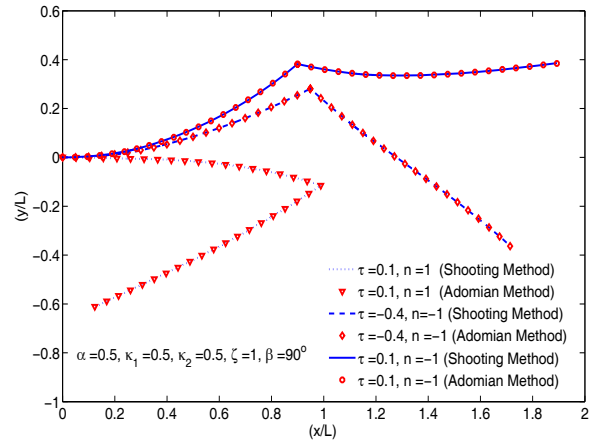


Figure 6b

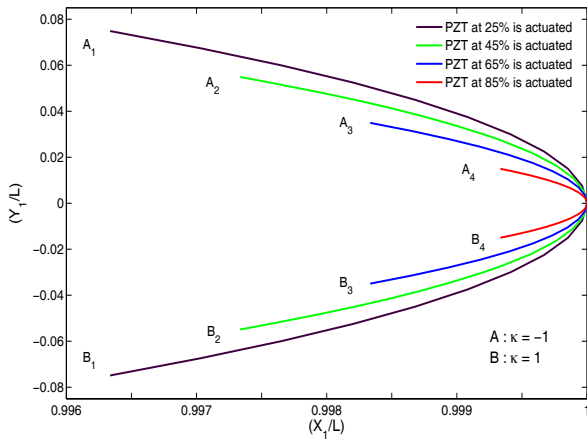


Figure 5b

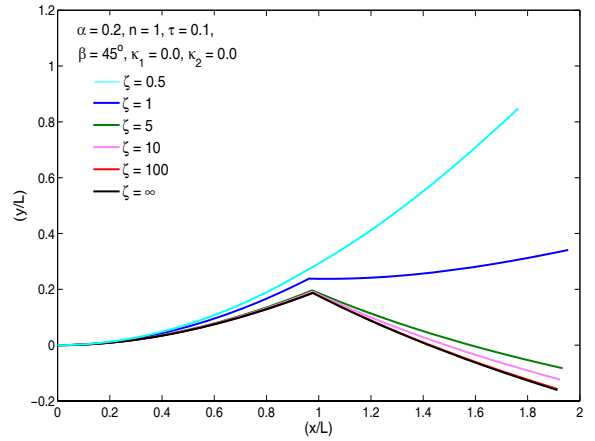


Figure 6c

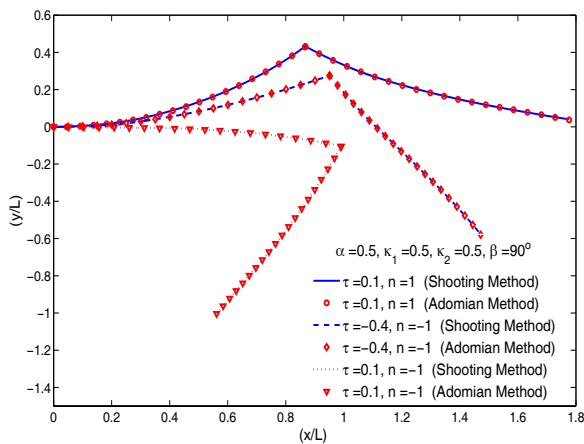


Figure 6a

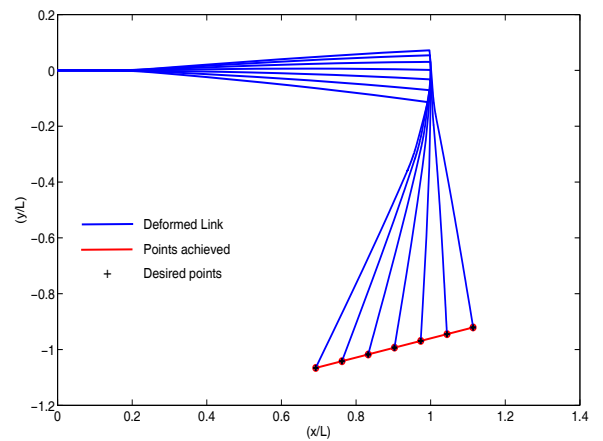


Figure 7a

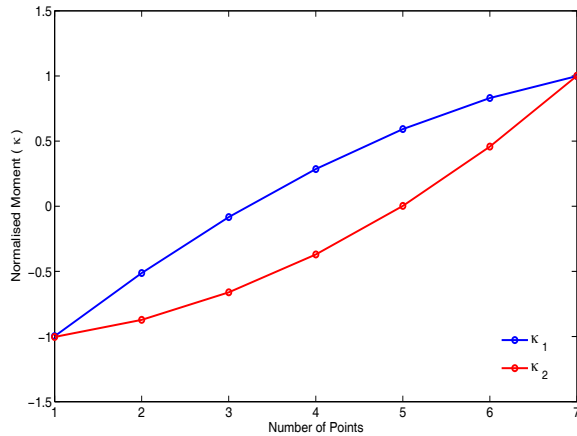


Figure 7b

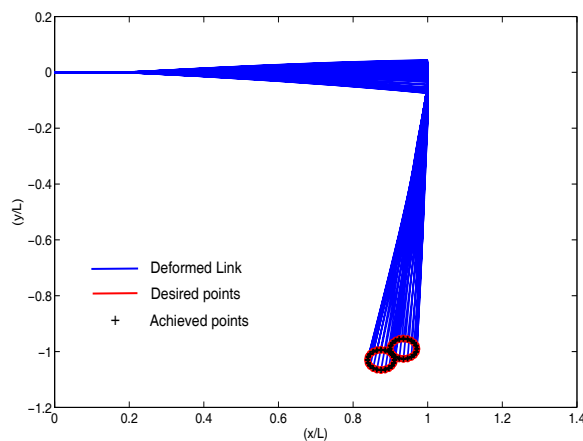


Figure 7c

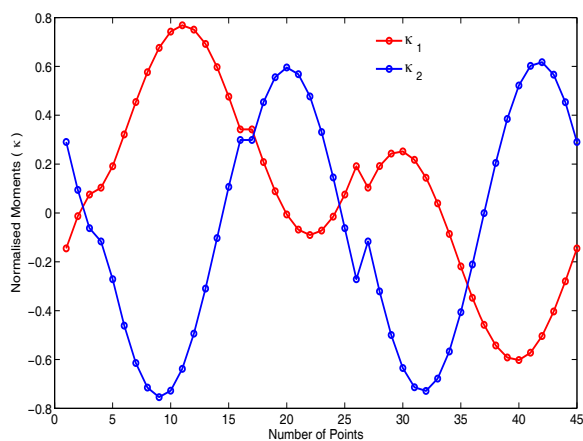


Figure 7d

Conclusions

A forward and an inverse analysis of path generation by a two link flexible mechanism actuated through piezoelectric actuators are presented. Two different methods viz., the non-linear shooting and Adomian decomposition have

been used. The methods are proved to be versatile for a range of loading parameters which may or may not cause inflection points within a link. Numerical results are included with different trajectories, viz., a straight line and a figure of eight.

In future these algorithms will be extended for forward and inverse analyses of closed chain fully compliant mechanisms.

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